## A Toy Model Study of Geometric Effects on Topological Transitions in Heavy-Ion Collisions

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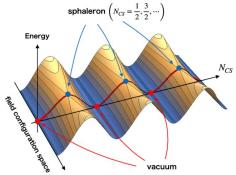
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#### Plan

- Motivation
- Description of toy model
- Toy model estimation of chiral chemical potential

#### Motivation

- QCD vacuum has non-trivial topological structure gauge field configurations like instantons and sphalerons.
- These transitions change the Chern–Simons number  $(N_{CS})$ , generating axial charge:  $\Delta Q_5 = 2N_f \Delta N_{CS}$ .
- Understanding the rate of topological transitions is crucial to interpret chirality-related observables in experiments. For example, Local chirality imbalance in presence of strong magnetic field (B) induces the Chiral Magnetic Effect (CME). Image (below) credit: @Phys. Rev. D 101, 096014



#### Motivation: Geometric evolution in HIC

- Before collision: two Lorentz-contracted 'pancake-shaped' nuclei approach each other.
- After collision: hot overlapping zone expands rapidly, becoming approximately spherical.
- This geometric evolution modifies:
  - Expansion rate
  - Cooling rate and lifetime of the medium
  - Temperature profile T( au) which controls transition rates like  $\Gamma_s \propto \alpha_s^5 T^4$

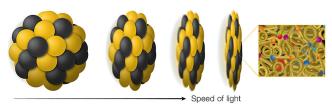


Image Credit: BNL

## Motivation: How geometric evolution influences topological transition rates

- Geometry determines expansion and thermalisation dynamics.
- Expansion affects temperature and lifetime both key factors for topological transitions.
- Therefore, geometric evolution may indirectly control the probability of exciting non-trivial topological configurations.

#### Key idea

Connect geometry-driven expansion to total integrated sphaleron transitions in a toy model.

## Toy model: setup and assumptions

#### Assumptions:

- Boost-invariant longitudinal expansion; transverse expansion included approximately.
- Temperature profile:  $T(\tau) = T_0(\tau_0/\tau)^{\alpha}$  ( $\alpha = \frac{1}{3}$  for Bjorken 1D).
- Initial energy density (Bjorken):  $\epsilon_0 = (dE_T/dy)/(A\tau_0)$ .
- EoS (toy):  $\epsilon = a_{eff}T^4 \rightarrow T_0 = (\epsilon_0/a_{eff})^{1/4}$ .
- Sphaleron rate:  $\Gamma_{sph}(T) = \kappa \alpha_s^5 T^4$ .

# Toy model: Analytic expression for integrated sphaleron count

For sphaleron count we start with:

$$N_{sph} = \int d^4x \, \Gamma_{sph}(T) \approx A \int_{\tau_0}^{\tau_f} d\tau \, \tau \, \Gamma_{sph}(T(\tau))$$

• For  $\Gamma_{sph}$  and  $T(\tau)$  we use:

$$\Gamma_{sph} \propto T^4$$

With  $T(\tau) = T_0(\tau_0/\tau)^{1/3}$ , we express  $\Gamma_{sph}$  as:

$$\tau \Gamma_{sph} \propto \tau^{-4/3}$$
.

 Integrated number of transitions in one unit rapidity (assuming boost invariance over rapidity chunk) for a transverse area A:

$$N_{sph} = A \kappa \alpha_s^5 T_0^5 \tau_0^{4/3} \frac{3}{2} \int_{\tau_0}^{\tau_f} d\tau \tau^{-1/3}.$$

## Toy model: Analytic expression for integrated sphaleron count

• To see geometry dependence explicitly, we substitute

$$T_0^4 = \epsilon_0/a_{eff} = \frac{1}{a_{eff}} \frac{dE_T/dy}{A\tau_0}$$
.

Finally:

$$N_{sph} = \kappa \,\alpha_s^5 \, \frac{dE_T/dy}{a_{eff}} \, \frac{3}{2} \, \tau_0^{1/3} \, \left(\tau_f^{2/3} - \tau_0^{2/3}\right).$$

#### Note:

- Smaller  $\tau_0$  translates into larger initial  $T_0$  (due to  $\epsilon_0 \propto 1/\tau_0$ ).
- $au_f$  increases with smaller  $au_0$  because  $T_0$  is larger  $\left( au_f \propto au_0 (T_0/T_f)^3\right)$ .
- The analytic expression shows competing trends:  $\tau_0$  growth vs reduced integration lower limit.

## Toy model estimation of Chiral chemical potential

Each topological transition produces a net axial charge:

$$\Delta Q_5 = 2 N_f \, \Delta N_{CS}$$

Here,  $\Delta Q_5$  is the net axial charge and  $\Delta N_{CS}$  is the change in the CHern-Simon number.

 The simplest kinetic equation (ignoring diffusion) involving axial charge density:

#### Kinetic equation

$$\frac{dn_5}{dt} = 2N_f \, \Gamma sph(T(t)) - \Gamma_f \, n_5(t)$$

 $n_5 =$ axial charge density

 $\Gamma_f$  = Chirality flipping (relaxation) rate.

## Toy model estimation of Chiral chemical potential

## Relation between axial charge density and chiral chemical potential

For small  $\mu_5$  (linear response) the axial charge density is related to chiral chemical potential by the axial susceptibility  $\chi_5$ :

$$n_5(t) = \chi_5(t)\mu_5(t)$$

#### Kinetic equation

Using the relation between axial charge and chiral chemical potential, the kinetic can be expressed in the following form:

$$\chi_5 \frac{d\mu_5}{dt} + \frac{d\chi_5}{dt} \mu_5 = 2N_f \Gamma_{sph} - \Gamma_f \chi_5 \mu_5$$

## Toy model estimation of Chiral chemical potential

• Under adiabatic approximation (ignoring the time derivative of  $\chi_5$ ) can be cast into the following form:

#### Kinetic equation under adiabatic approximation

$$\frac{d\mu_5}{dt} = \frac{2N_f \, \Gamma_{sph}}{\chi_5} - \Gamma_f \, \mu_5$$

Upon integration, we can have an estimation of the chiral chemical potential.

Under steady state approximation:

$$\frac{d\mu_5}{dt} = 0$$

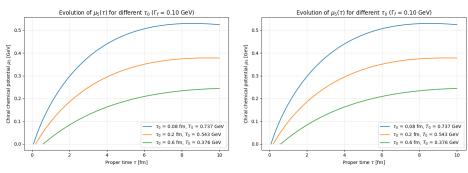
#### Chiral chemical potential under steady state approximation

Under steady state approximation, the chiral chemical potential assumes the following simpler form:

$$\mu_5^{ss}(T) \approx \frac{2N_f\;\Gamma_{sph}}{\chi_5\;\Gamma_f}$$

### Toy model estimation of chiral chemical potential

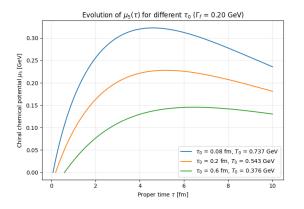
#### $\mu_5$ evolution with example $\Gamma_f$ values



- Initial rise:  $\mu_5$  increases rapidly when the temperature is highest and sphaleron transition rate term dominates
- Saturation region: As n<sub>5</sub> builds up, the damping term balances the sphalerons source term. The system achieves a quasi-steady state.

### Toy model estimation of chiral chemical potential

 $\mu_5$  evolution with example  $\Gamma_f$  values



• Decay phase: As the QGP cools down both  $\Gamma_{sph}$  and  $\chi_5$  decreases sharply and the damping term dominates and give rise to slow decay of  $\mu_5$ .

Thank you

Thank You!

## Supplementary