The surprising agreement of hypertriton yields in high energy HICs with thermal model

Hypertriton puzzle in relativistic heavy-ion collisions

Thomas Cohen, Maneesha Pradeep, Phys. Rev. C 111 (2025) 5, 054917

University of Maryland, College Park & Indian Institute of Science, Bangalore

India-JINR Workshop 2025 at NISER, India

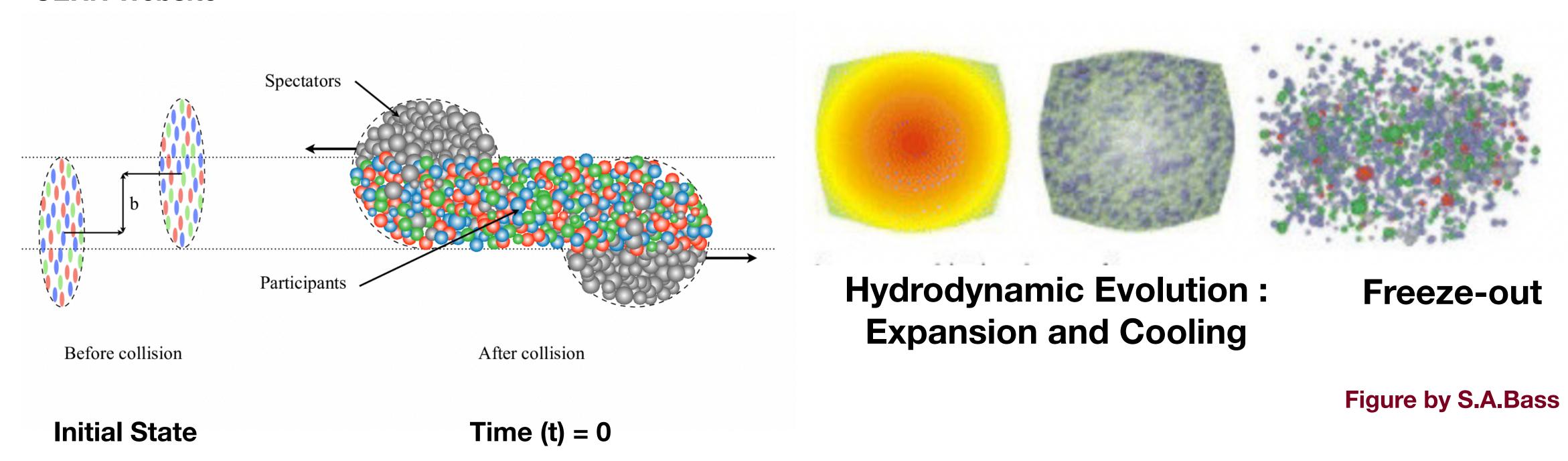






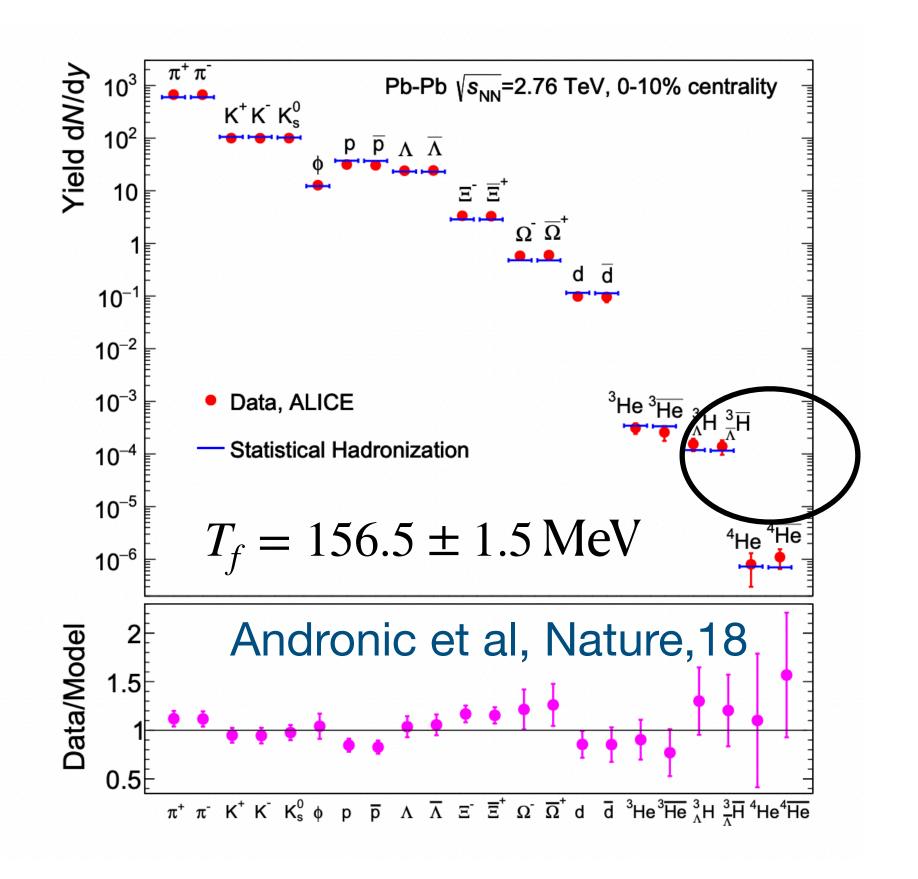
Heavy-Ion Collisions

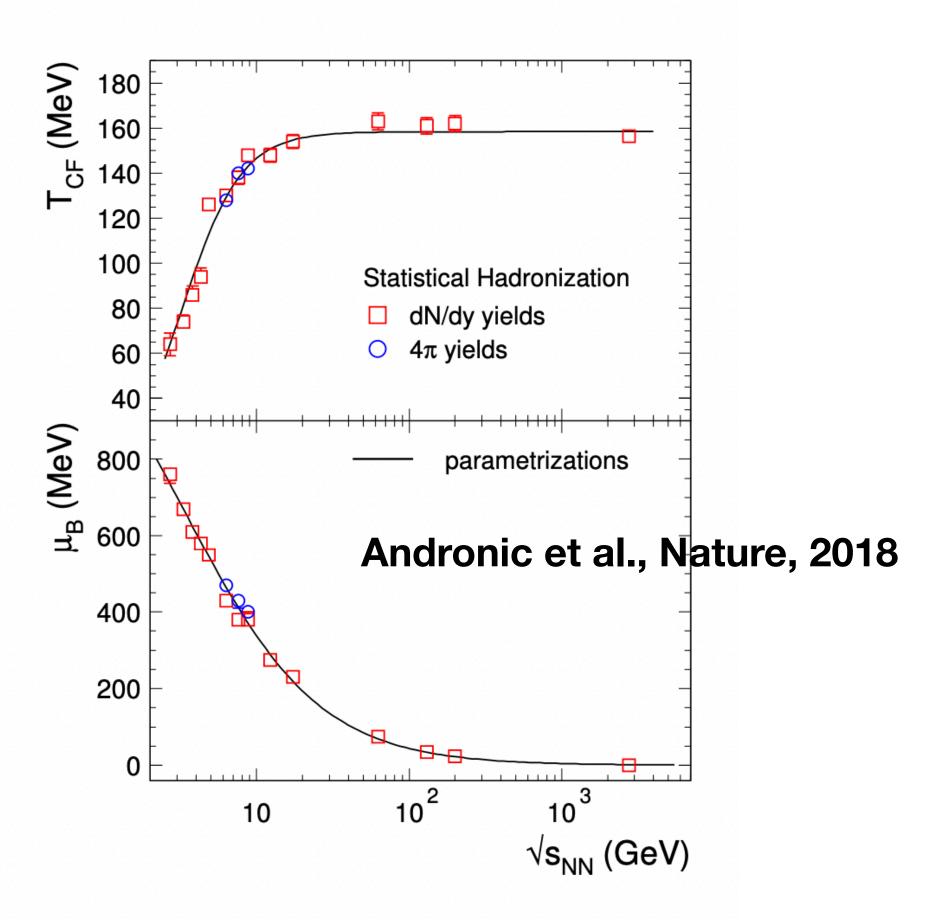
CERN Website



- \circ QGP thermalizes in a relatively short amount of time \sim 1fm/c \sim 10^{-23} s
- Freeze-out at about 10 fm/c ~10^{-22}s
- The event by event distribution of the particle multiplicities are frozen at freeze-out and are measured at the detectors.

Statistical Hadronization Model at freeze-out



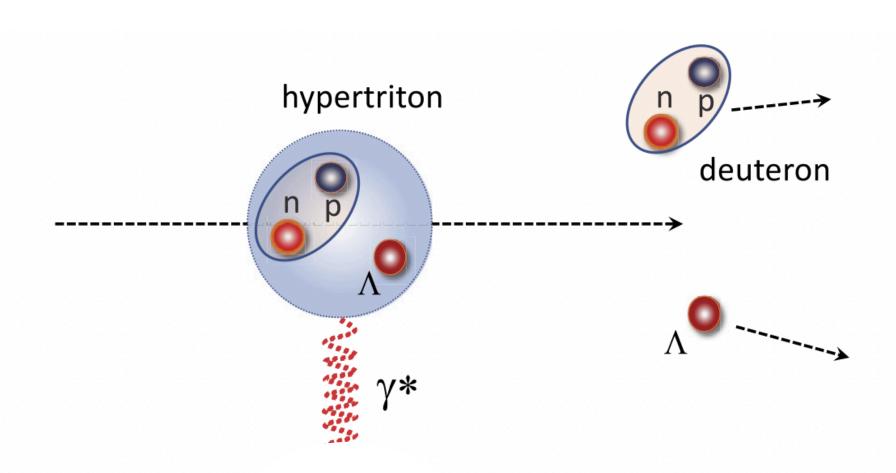


- Mid-rapidity yields consistent with a Statistical Hadron Resonance Gas model across a broad range of collision energies
- Model assumes: chemical freeze-out → yields fixed by an ideal gas of hadrons & nuclei (free-space masses) in equilibrium.
- Stunning phenomenological success: spans nine orders of magnitude in abundances including light (hyper) nuclei.

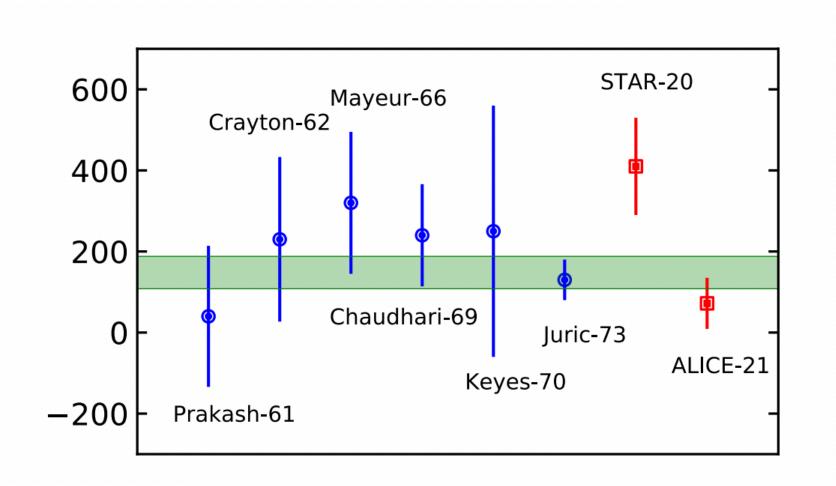
Physical picture at freeze-out

- There is a gas of almost non-interacting hadrons and light nuclei in thermal equilibrium at 156 MeV at freeze-out within the confines of the fireball boundary
- Particles are well localized
- Binding energy of light nuclei, typically few MeV or keV << 156 MeV
- Still plausible, if the gas is dilute enough

Hypertritons in HICs-"Snowballs in hell"



B.E $\sim 148 \pm 40 \,\text{keV} \ll 156 \,\text{MeV}$



Bertulani, 22

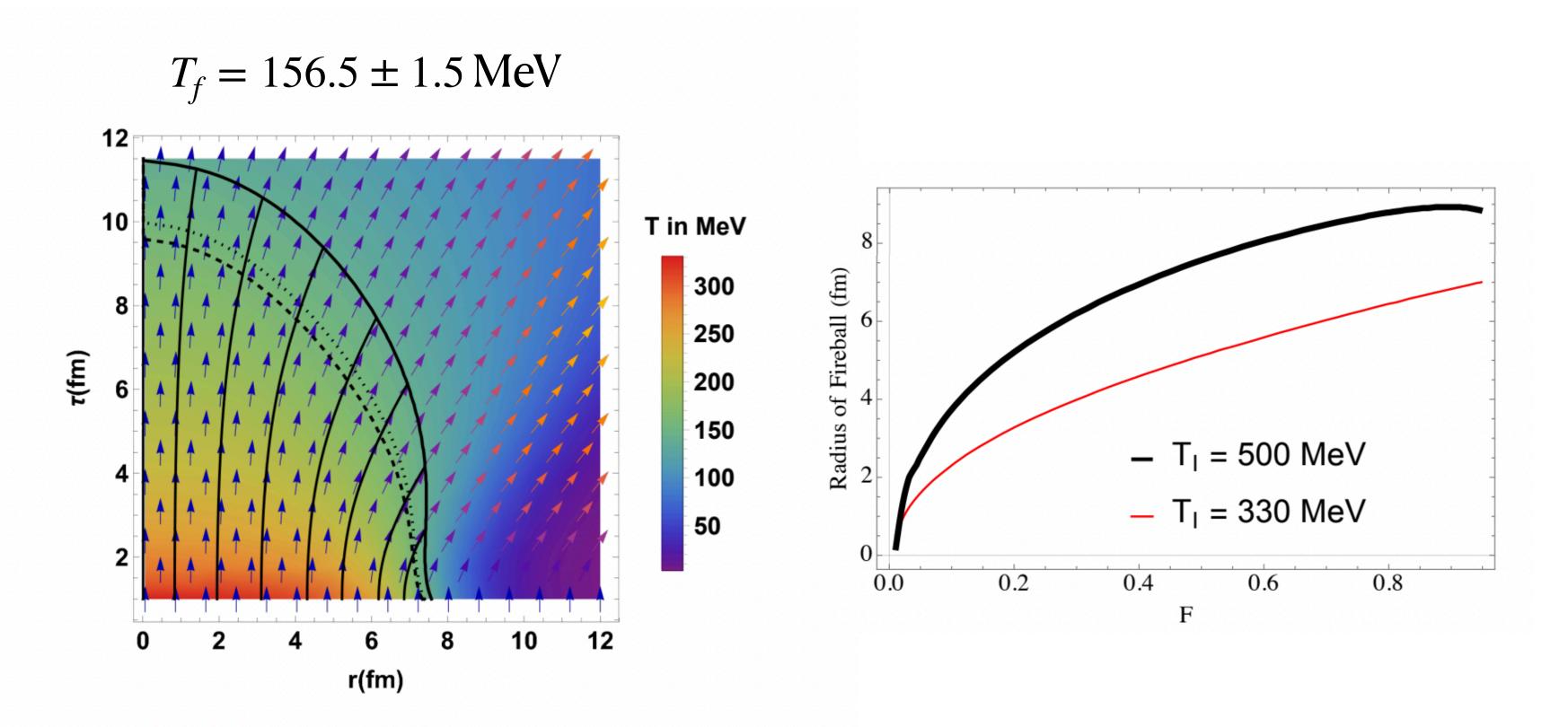
- Weakly bound state of deuteron and Lambda baryon
- A single thermal pion can dissociate H_{Λ} by quasi-free scattering / resonant process ~ short lifetimes of a few fm
- Root mean-squared radius $\sim 10\,\mathrm{fm}$ ~ Size of the fireball?

Quantum mechanical hypertriton comes in all sizes!

$$V(r) = \begin{cases} -V_0, & r < a, \\ 0, & r \ge a, \end{cases} \xrightarrow{\begin{bmatrix} 0.30 \\ 0.25 \\ 0.05 \\ 0.00 \end{bmatrix}} \xrightarrow{\begin{array}{c} 0.25 \\ 0.15 \\ 0.00 \\ 0 \end{array}} \xrightarrow{\begin{array}{c} 0.25 \\ 0.10 \\ 0.05 \\ 0.00 \\ 0 \end{array}} \xrightarrow{\begin{array}{c} 0.25 \\ 0.10 \\ 0.05 \\ 0.00 \\ 0 \end{array}} \xrightarrow{\begin{array}{c} 0.25 \\ 0.2 \\ 0.00 \\ 0 \end{array}} \xrightarrow{\begin{array}{c} 0.4 \\ 0.2 \\ 0.00 \\ 0 \end{array}} \xrightarrow{\begin{array}{c} 0.4 \\ 0.2 \\ 0.00 \\ 0 \end{array}} \xrightarrow{\begin{array}{c} 0.4 \\ 0.2 \\ 0.00 \\ 0 \end{array}} \xrightarrow{\begin{array}{c} 0.4 \\ 0.2 \\ 0.00 \\ 0 \end{array}} \xrightarrow{\begin{array}{c} 0.4 \\ 0.2 \\ 0.00 \\ 0 \end{array}} \xrightarrow{\begin{array}{c} 0.4 \\ 0.2 \\ 0.00 \\ 0 \end{array}} \xrightarrow{\begin{array}{c} 0.4 \\ 0.2 \\ 0.00 \\ 0 \end{array}} \xrightarrow{\begin{array}{c} 0.4 \\ 0.2 \\ 0.00 \\ 0 \end{array}} \xrightarrow{\begin{array}{c} 0.4 \\ 0.2 \\ 0.00 \\ 0 \end{array}} \xrightarrow{\begin{array}{c} 0.4 \\ 0.2 \\ 0.00 \\ 0 \end{array}} \xrightarrow{\begin{array}{c} 0.4 \\ 0.2 \\ 0.00 \\ 0 \end{array}} \xrightarrow{\begin{array}{c} 0.4 \\ 0.2 \\ 0.00 \\ 0 \end{array}} \xrightarrow{\begin{array}{c} 0.4 \\ 0.2 \\ 0.00 \\ 0 \end{array}} \xrightarrow{\begin{array}{c} 0.4 \\ 0.2 \\ 0.00 \\ 0 \end{array}} \xrightarrow{\begin{array}{c} 0.4 \\ 0.2 \\ 0.00 \\ 0 \end{array}} \xrightarrow{\begin{array}{c} 0.4 \\ 0.2 \\ 0.00 \\ 0 \end{array}} \xrightarrow{\begin{array}{c} 0.4 \\ 0.2 \\ 0.00 \\ 0 \end{array}} \xrightarrow{\begin{array}{c} 0.4 \\ 0.2 \\ 0.00 \\ 0 \end{array}} \xrightarrow{\begin{array}{c} 0.4 \\ 0.2 \\ 0.00 \\ 0 \end{array}} \xrightarrow{\begin{array}{c} 0.4 \\ 0.2 \\ 0.00 \\ 0 \end{array}} \xrightarrow{\begin{array}{c} 0.4 \\ 0.2 \\ 0.00 \\ 0 \end{array}} \xrightarrow{\begin{array}{c} 0.4 \\ 0.2 \\ 0.00 \\ 0 \end{array}} \xrightarrow{\begin{array}{c} 0.4 \\ 0.2 \\ 0.00 \\ 0 \end{array}} \xrightarrow{\begin{array}{c} 0.4 \\ 0.2 \\ 0.00 \\ 0 \end{array}} \xrightarrow{\begin{array}{c} 0.4 \\ 0.2 \\ 0.00 \\ 0 \end{array}} \xrightarrow{\begin{array}{c} 0.4 \\ 0.2 \\ 0.00 \\ 0 \end{array}} \xrightarrow{\begin{array}{c} 0.4 \\ 0.2 \\ 0.00 \\ 0 \end{array}} \xrightarrow{\begin{array}{c} 0.4 \\ 0.2 \\ 0.00 \\ 0 \end{array}} \xrightarrow{\begin{array}{c} 0.4 \\ 0.2 \\ 0.00 \\ 0 \end{array}} \xrightarrow{\begin{array}{c} 0.4 \\ 0.2 \\ 0.00 \\ 0 \end{array}} \xrightarrow{\begin{array}{c} 0.4 \\ 0.2 \\ 0.00 \\ 0 \end{array}} \xrightarrow{\begin{array}{c} 0.4 \\ 0.2 \\ 0.00 \\ 0 \end{array}} \xrightarrow{\begin{array}{c} 0.4 \\ 0.2 \\ 0.00 \\ 0 \end{array}} \xrightarrow{\begin{array}{c} 0.4 \\ 0.2 \\ 0.00 \\ 0 \end{array}} \xrightarrow{\begin{array}{c} 0.4 \\ 0.2 \\ 0.00 \\ 0 \end{array}} \xrightarrow{\begin{array}{c} 0.4 \\ 0.2 \\ 0.00 \\ 0 \end{array}} \xrightarrow{\begin{array}{c} 0.4 \\ 0.2 \\ 0.00 \\ 0 \end{array}} \xrightarrow{\begin{array}{c} 0.4 \\ 0.2 \\ 0.00 \\ 0 \end{array}} \xrightarrow{\begin{array}{c} 0.4 \\ 0.2 \\ 0.00 \\ 0 \end{array}} \xrightarrow{\begin{array}{c} 0.4 \\ 0.2 \\ 0.00 \\ 0 \end{array}} \xrightarrow{\begin{array}{c} 0.4 \\ 0.2 \\ 0.00 \\ 0 \end{array}} \xrightarrow{\begin{array}{c} 0.4 \\ 0.2 \\ 0.00 \\ 0 \end{array}} \xrightarrow{\begin{array}{c} 0.4 \\ 0.2 \\ 0.00 \\ 0 \end{array}} \xrightarrow{\begin{array}{c} 0.4 \\ 0.2 \\ 0.00 \\ 0 \end{array}} \xrightarrow{\begin{array}{c} 0.4 \\ 0.2 \\ 0.00 \\ 0 \end{array}} \xrightarrow{\begin{array}{c} 0.4 \\ 0.2 \\ 0.00 \\ 0 \end{array}} \xrightarrow{\begin{array}{c} 0.4 \\ 0.2 \\ 0.00 \\ 0.0$$

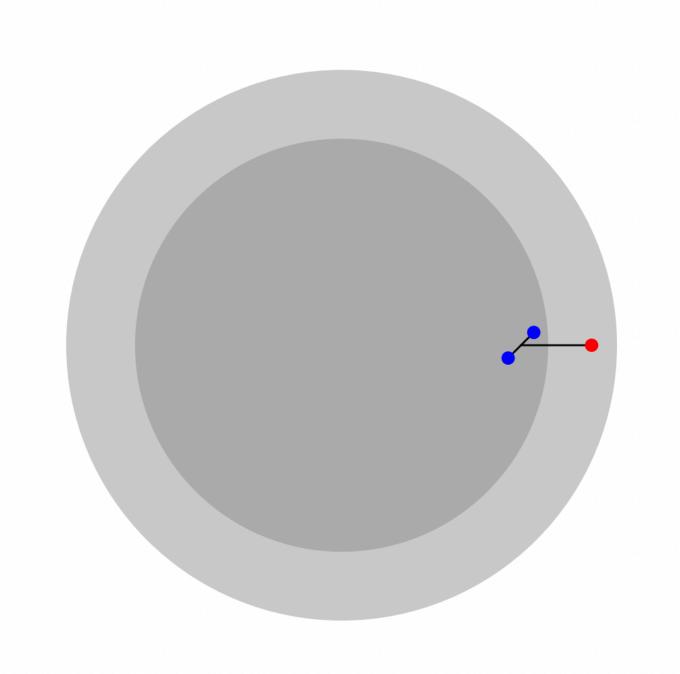
- Probability distribution for the radial separation between Λ and d has a long ranged tail
- Magnitude of the short distance part of the H_{Λ} wave function is non-negligible, even though small due to large binding energy
- Deuteron itself has a root mean-squared radius of 8.1 fm and a weak B.E of 2.24 MeV
- We model both H_{Λ} and d by finite square well potential specified by the depth of the potential and finite radial extent

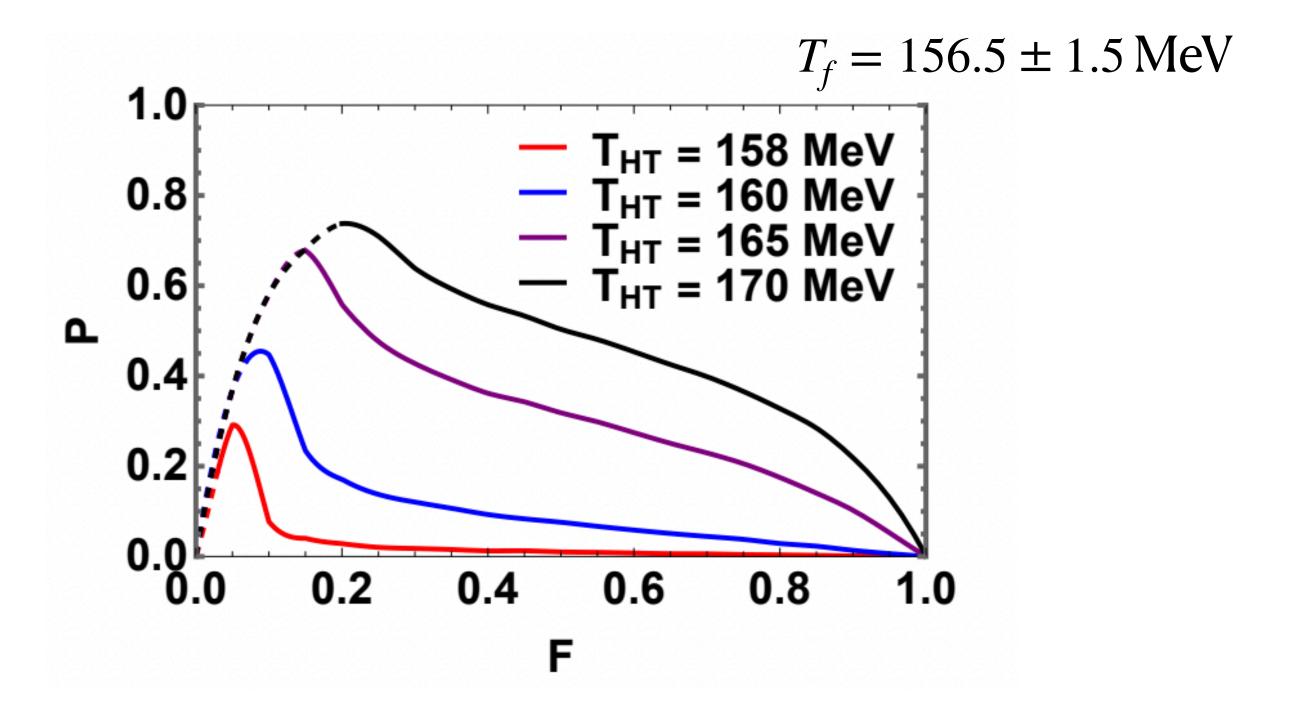
Fireball radius decreases as the system cools with time



• The radius of the remaining fireball decreases as the fraction of energy density remaining in the fireball decreases

Will the hypertriton fit inside the fireball?



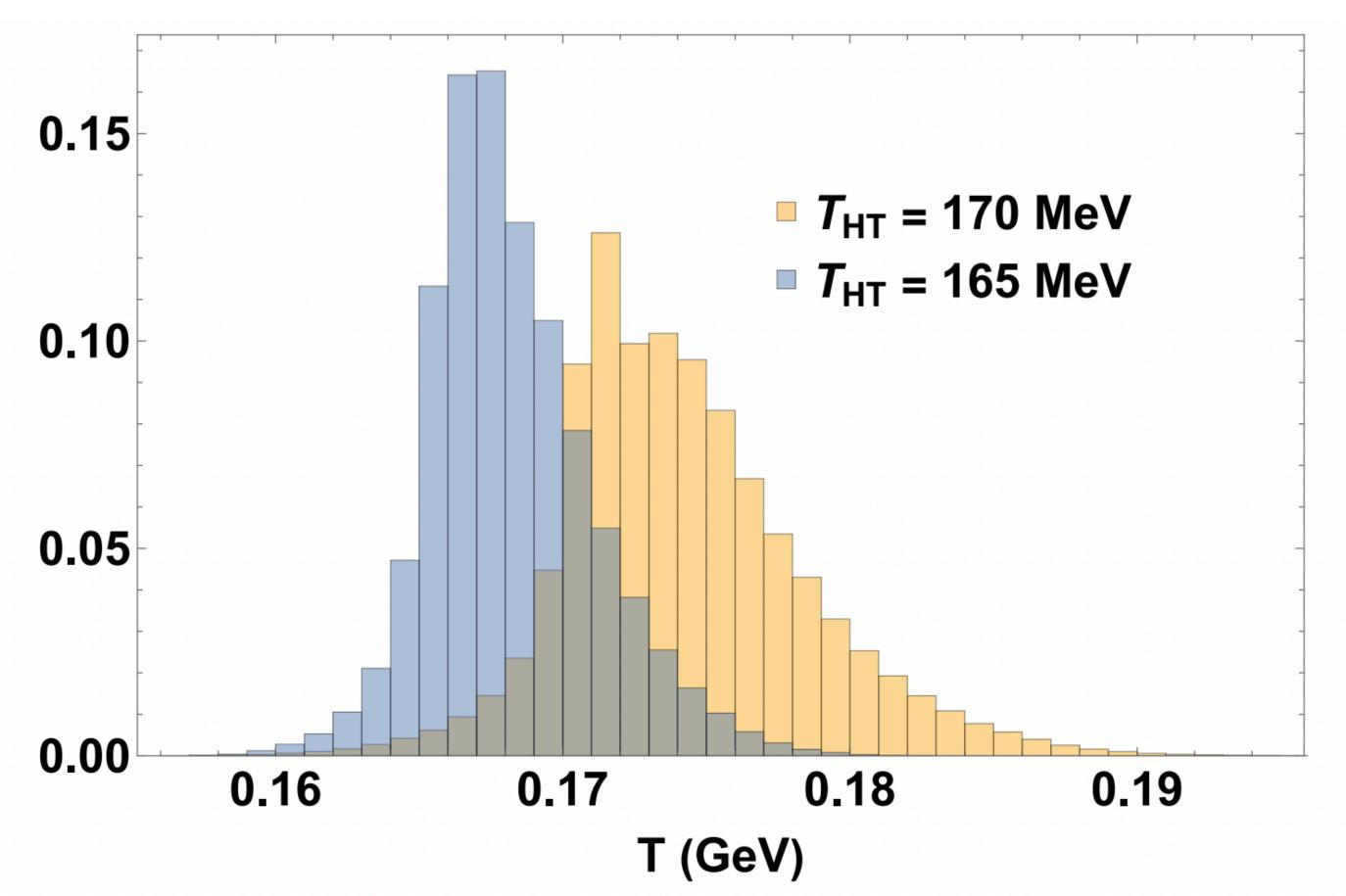


Suppose the center of mass of the H_{Λ} is at a temperature $T_{\rm HT}$, what is the probability of the proton, neutron and Λ being inside the fireball when fraction of the initial energy density, F is remaining?

$$P_{\rm IN} < 30\%$$
 for $T_{\rm HT} = 170$ MeV, $P_{\rm IN} < 10\%$ for $T_{\rm HT} = 160$ MeV

Temperature gradients across the size of hypertriton

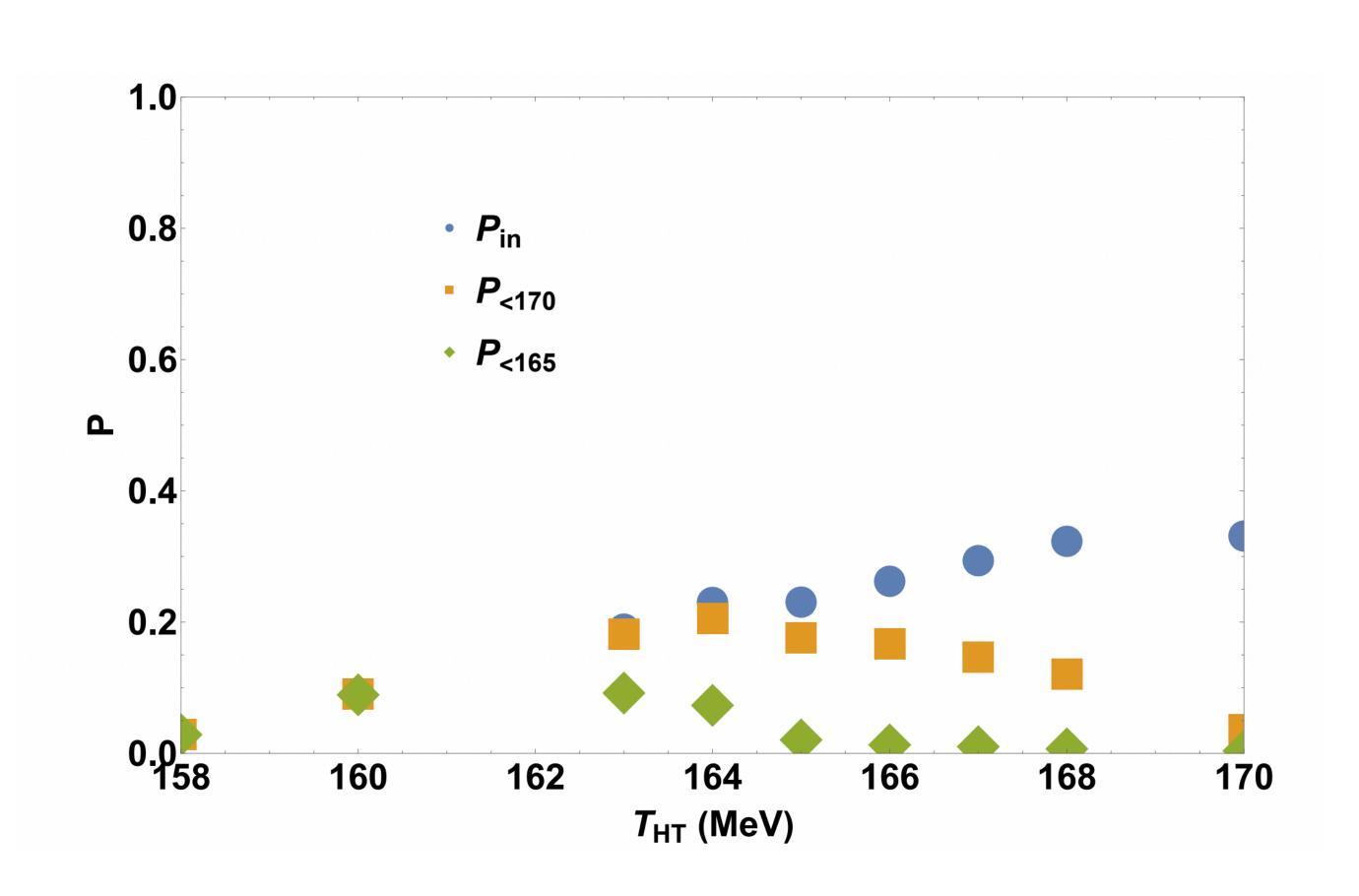
Distribution of maximum of the temperature at the location of its constituents



$$T_f = 156.5 \pm 1.5 \,\text{MeV}$$

- Temperature variations across the size of the hypertriton is huge
- $R_T = \langle N \rangle_T / \langle N \rangle_{156.5}$
- $R_{170} = 5.2$
- $R_{165} = 2.9$
- $R_{160} = 1.6$

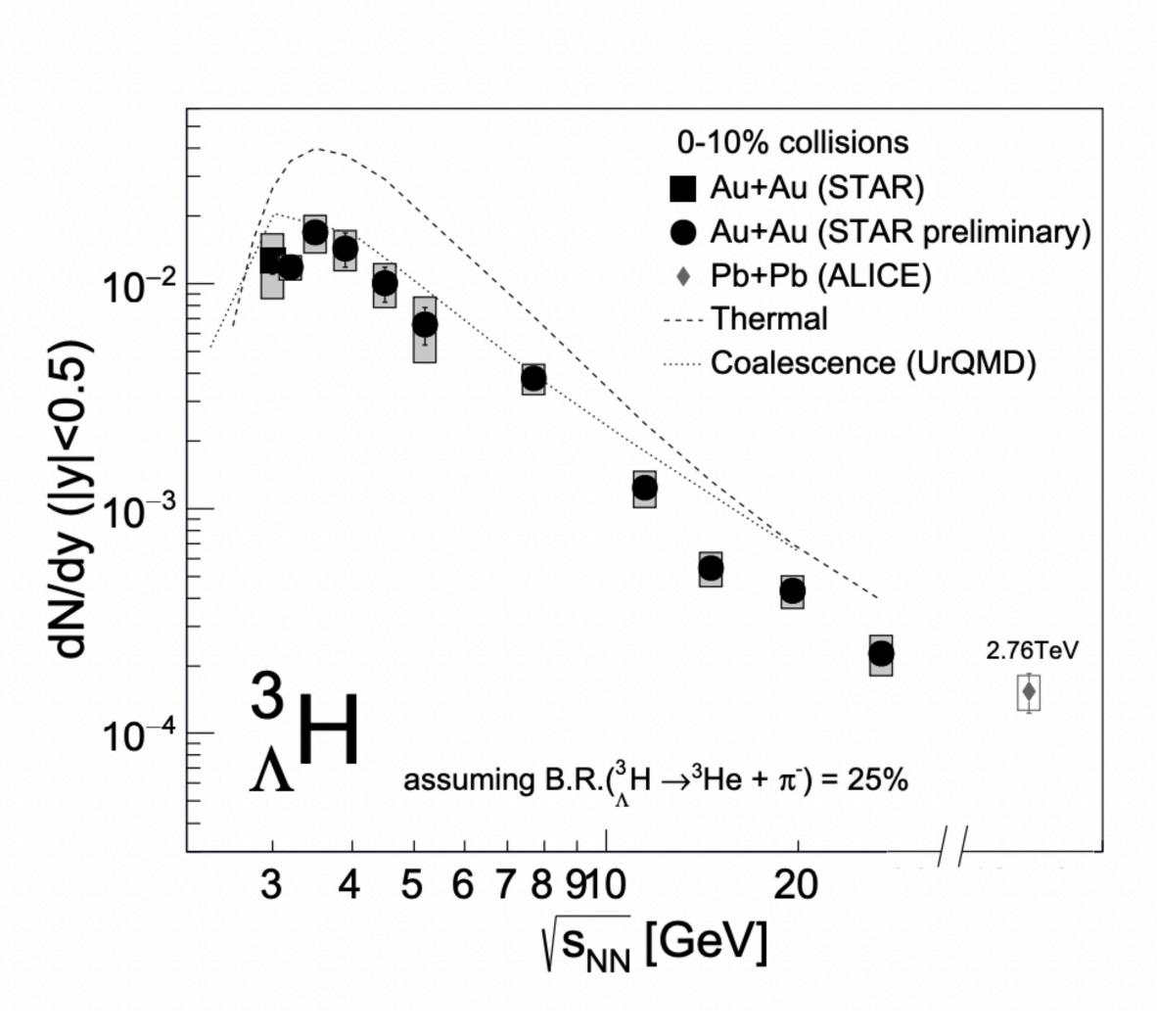
The hypertriton puzzle



$$T_f = 156.5 \pm 1.5 \,\text{MeV}$$

Consistency with data and quantum mechanics demands the hypertriton reside in regions below 160 MeV, where the likelihood of all constituents fitting is under 10%.

Thermal agreement of hypertriton production as a function of collision energy



- Discrepancies between experiment and thermal model increase at lower energies
- Coalescence models more parameters and features can also explain
- Hypertriton production is maximized at energies which can probed at NICA
- Event by event fluctuation measurements with large statistics

Summary

- The weakly bound hypertriton demands a quantum treatment on length scales comparable to the fireball at freeze-out.
- A simple quantum model yields a conservative ≤10% bound on the probability that a thermally produced hypertriton stays inside the freeze-out shell.
- Refinements to the quantum mechanical descriptions is likely to only decrease this estimate as binding energy decreases in medium.
- Hypertriton production remains a puzzle, and NICA offers an opportunity to investigate it further.

Thank you!

Additional Slides

Hypertriton inside a fireball

