

# Series expansion of hypergeometric functions about their parameters using MultiHypExp

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and

Comput.Phys.Commun. 297 (2024) 109060 (arXiv:2306.11718)

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*Advanced School & Workshop on Multiloop Scattering Amplitudes  
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# Outline

Motivation

Definitions

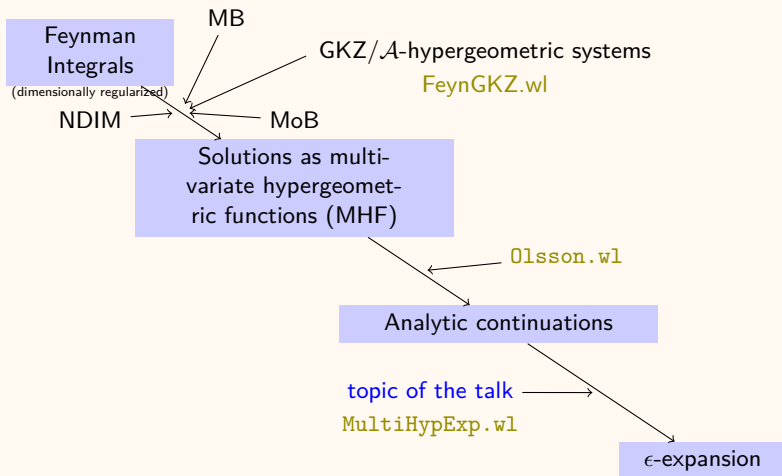
Feynman integrals and hypergeometric functions

Series expansion

Algorithm

Mathematica package : MultiHypExp

# The Big Picture



## Definitions

► **Pochhammer symbol :**

$$\begin{aligned}(x)_n &= \frac{\Gamma(x+n)}{\Gamma(x)}, \\ &= x(x+1)\dots(x+n-1), \quad x \in \mathbb{C} \setminus \mathbb{Z}_0^-, n \in \mathbb{Z}_0^+ \\ (1)_n &= n!\end{aligned}$$

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- ▶ General hypergeometric functions

$${}_pF_{p-1}(a_1, \dots, a_p; b_1, \dots, b_q; z) = \sum_{n=0}^{\infty} \frac{(a_1)_n \cdots (a_p)_n}{(b_1)_n \cdots (b_{p-1})_n} \frac{z^n}{n!}, \quad |z| < 1$$

## Definitions

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$$\text{Appell } F_2(a, b_1, b_2; c_1, c_2; x, y) = \sum_{m,n=0}^{\infty} \frac{(a)_{m+n} (b_1)_m (b_2)_n}{(c_1)_m (c_2)_n} \frac{x^m y^n}{m!n!}$$

valid for  $|x| + |y| < 1$

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- ▶ **Lauricella functions:**  $F_A, F_B, F_C$  and  $F_D$

$$F_C^{(3)} = \sum_{n_1, n_2, n_3=0}^{\infty} \frac{(a_1)_{n_1+n_2+n_3} (a_2)_{n_1+n_2+n_3}}{(c_1)_{n_1} (c_2)_{n_3} (c_3)_{n_2}} \frac{z_1^{n_1} z_2^{n_2} z_3^{n_3}}{n_1! n_2! n_3!}$$

with domain of convergence :  $\sqrt{|z_1|} + \sqrt{|z_2|} + \sqrt{|z_3|} < 1$

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- ▶ One loop three-point function :

$$F_2(\epsilon + 1, 1, 1; \epsilon + 1, 2 - \epsilon; x, y), \\ F_2(1, 1 - \epsilon, 1; 1 - \epsilon, 2 - \epsilon; x, y), \dots$$

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- ▶ The sunset integral with unequal masses : Berends et. al. [2]

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- ▶  $\epsilon$ -expansion of the multivariate hypergeometric functions (MHF) are needed

## From the literature

- ▶ Each of the representations of the MHF can be used
  - ▶ Series
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- ▶ Available packages:
  - ▶ **Analytical** :
    - ▶ HypExp , HypExp2 (T. Huber et. al. [3, 4])
    - ▶ XSummer (S. Moch et. al. [5])
    - ▶ nestedsums (S. Weinzierl [6, 7])
    - ▶ RISC packages : Sigma (C. Schneider [8]) , EvaluateMultiSums (C. Schneider [9]) , HarmonicSums (J. Ablinger [10])
  - ▶ **Numerical** : NumExp (Z.W. Huang et. al. [11])

## Example 0

► Case I :

$${}_2F_1(1, 1; \epsilon + 1; x) = \sum_{m=0}^{\infty} \frac{(1)_m (1)_m}{(\epsilon + 1)_m} \frac{x^m}{m!} = \sum_{m=0}^{\infty} x^m + O(\epsilon) = \frac{1}{1-x} + O(\epsilon)$$

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- Case II :

$${}_2F_1(1, 1; \epsilon - 1; x) = \sum_{m=0}^{\infty} \frac{(1)_m(1)_m}{(-1 + \epsilon)_m} \frac{x^m}{m!} = ?$$

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- ▶ Proposed in [12]

$$\begin{aligned} {}_2F_1(1, 1; \epsilon - 1; x) &= H(\epsilon) \bullet {}_2F_1(1, 1; \epsilon + 1; x) \\ &= \left[ \frac{1}{\epsilon} \left[ \frac{x}{(x-1)} + \frac{3x-1}{(x-1)} x \partial_x \right] + O(\epsilon^0) \right] \bullet \left[ \frac{1}{1-x} + O(\epsilon) \right] \\ &= \frac{1}{\epsilon} \frac{2x^2}{(x-1)^3} + O(\epsilon^0) \end{aligned}$$

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- ▶ MHF with *singular* parameters may have Laurent series expansion

## Algorithm

- ▶ **Step 1:** Check if the Pochhammer parameters of the given function ( $F(\epsilon)$ ) are *singular* or not
- ▶ **Step 2:** If those are non-singular, find the Taylor expansion of  $F(\epsilon)$
- ▶ **Step 3:** If any of the Pochhammer parameter of  $F(\epsilon)$  is *singular* then find a new function ( $G(\epsilon)$ ) by replacing

*singular* Pochhammer  $\longrightarrow$  *non-singular* Pochhammer

- ▶ **Step 4:** Relate them using a differential operator ( $H$ )

$$F(\epsilon) = H(\epsilon) \bullet G(\epsilon)$$
$$\left[ \sum_{i=-n}^{\infty} \epsilon^i H_i \right] \bullet \left[ \sum_{j=0}^{\infty} \epsilon^j G_j \right]$$

## Step 1 & 3 : Checking the Pochhammers

► *Singular* Pochhammers :

1. When one or more lower Pochhammer parameters (i.e., Pochhammer parameters in the denominator) are of the form

$$(-p + q\epsilon)_p$$

2. When one or more upper Pochhammer parameters (i.e., Pochhammer parameters in the numerator) are of the form

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- The Gauss  ${}_2F_1$  example

$$A_1 := \sum_{m=0}^{\infty} \frac{(\epsilon)_m (-\epsilon)_m x^m}{(\epsilon - 1)_m m!} \quad , \quad A_2 := \sum_{m=0}^{\infty} \frac{(\epsilon)_m (-\epsilon)_m x^m}{(\epsilon + 1)_m m!}$$

*singular* Pochhammer in  $A_1$  :  $(-1 + \epsilon)_m$   
*non-singular* Pochhammer in  $A_2$  :  $(1 + \epsilon)_m$

## Step 2 : Taylor Expansion

Obtaining DE

- ▶ From the definition of Taylor expansion

$$F(\epsilon) = \sum_{i=0}^{\infty} \frac{\epsilon^i}{i!} \left. \frac{d^i}{d\epsilon^i} F(\epsilon) \right|_{\epsilon=0}$$

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Expansion around [integer parameters](#)

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### Expansion around **integer parameters**

- ▶ Consider

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$$\frac{A_{n+1}}{A_n} = \frac{(a+n)(b+n)}{(n+1)(c+n)} = \frac{g(n)}{h(n)}$$

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- ▶ The annihilator

$$L = \left[ h(\theta) \frac{1}{x} - g(\theta) \right]$$

where  $\theta = x\partial_x$



## Step 2 : Taylor Expansion

### Pfaff System

- ▶ For Gauss  ${}_2F_1$  :  $L \bullet {}_2F_1(a, b; c; x) = 0$

$$L = -ab + (c - x(a + b + 1))\partial_x - (x - 1)x\partial_x^2$$

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- ▶ Consider  $g = ({}_2F_1, x\partial_x \bullet {}_2F_1)^T$

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- ▶ The length of the vector  $g = \text{Holonomic rank of the system}$
- ▶ Find a transformation  $T$  to bring the system into *canonical* form (J. Henn [13])

$$dg' = \epsilon \Omega' g'$$

with

$$g = Tg' \quad , \quad \Omega' = T^{-1}\Omega T - T^{-1}dT$$

## Step 2 : Taylor Expansion

### Example

- For our example of  $A_2 = {}_2F_1(\epsilon, -\epsilon; 1 + \epsilon; x)$

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- In this case,

$$T = \begin{pmatrix} 1 & 0 \\ 0 & \epsilon \end{pmatrix}, \quad \Omega' = \begin{pmatrix} 0 & \frac{1}{x} \\ \frac{1}{x-1} & \frac{1}{x-1} - \frac{1}{x} \end{pmatrix}$$

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$$T = \begin{pmatrix} 1 & 0 \\ 0 & \epsilon \end{pmatrix}, \quad \Omega' = \begin{pmatrix} 0 & \frac{1}{x} \\ \frac{1}{x-1} & \frac{1}{x-1} - \frac{1}{x} \end{pmatrix}$$

- ▶ **Boundary Condition** : At  $x = 0$

$$g = (1, 0)^T$$



## Step 2 : Taylor Expansion

### Example

- ▶ For our example of  $A_2 = {}_2F_1(\epsilon, -\epsilon; 1 + \epsilon; x)$

$$\Omega = \begin{pmatrix} 0 & \frac{1}{x} \\ \frac{\epsilon^2}{x-1} & \frac{\epsilon}{x-1} - \frac{\epsilon}{x} \end{pmatrix}$$

- ▶ In this case,

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- ▶ **Boundary Condition** : At  $x = 0$

$$g = (1, 0)^T$$

- ▶ **Solution** :

$$A_2 = 1 + \epsilon^2 G(0, 1; x) + \epsilon^3 (-G(0, 0, 1; x) + G(0, 1, 1; x)) + O(\epsilon^4)$$

## Step 4 : Differential operator

### Contiguous relations

There exist contiguous relations that relate

$${}_2F_1(a \pm 1, b; c; x) \quad , \quad {}_2F_1(a, b \pm 1; c; x) \quad , \quad {}_2F_1(a, b; c \pm 1; x)$$

These can be obtained by applying differential operators

► *Example:*

The unit step down operator for the Gauss hypergeometric series is given by  $H(c) = \frac{1}{c}(\theta + c)$ , i.e.,

$${}_2F_1(a, b; c; x) = H(c) \bullet {}_2F_1(a, b; c + 1; x)$$

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▶ Another example

$${}_2F_1(a + 1, b; c; x) = \frac{1}{a}(\theta + a) \bullet {}_2F_1(a, b; c; x)$$

## Step 4 : Differential operator

### Step Down Operators

- ▶ If needed, apply the step down operator multiple times

$${}_2F_1(a, b; c - 1; x) = H(c - 1)H(c) \bullet {}_2F_1(a, b; c + 1; x)$$

## Step 4 : Differential operator

### Step Down Operators

- ▶ If needed, apply the step down operator multiple times

$${}_2F_1(a, b; c - 1; x) = H(c - 1)H(c) \bullet {}_2F_1(a, b; c + 1; x)$$

- ▶ Take quotient by the annihilator  $L$

$$L \bullet {}_2F_1(a, b; c + 1, x) = 0$$

$$L = [-ab + (-x(a + b + 1) + c + 1)\partial_x - (x - 1)x\partial_x^2]$$

## Step 4 : Differential operator

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$$L = [-ab + (-x(a + b + 1) + c + 1)\partial_x - (x - 1)x\partial_x^2]$$

- ▶ The step down operator :

$$\begin{aligned} H &= H(c - 1)H(c) \\ &= \left(1 - \frac{abx}{(c - 1)c(x - 1)}\right) - \frac{x(a + b + 1) - 2cx + c - 1}{(c - 1)c(x - 1)}\theta \end{aligned}$$

## Step 4 : Differential operator

Example



$$A_1 := \sum_{m=0}^{\infty} \frac{(\epsilon)_m (-\epsilon)_m}{(\epsilon - 1)_m} \frac{x^m}{m!} \quad , \quad A_2 := \sum_{m=0}^{\infty} \frac{(\epsilon)_m (-\epsilon)_m}{(\epsilon + 1)_m} \frac{x^m}{m!}$$

## Step 4 : Differential operator

Example



$$A_1 := \sum_{m=0}^{\infty} \frac{(\epsilon)_m (-\epsilon)_m}{(\epsilon-1)_m} \frac{x^m}{m!} \quad , \quad A_2 := \sum_{m=0}^{\infty} \frac{(\epsilon)_m (-\epsilon)_m}{(\epsilon+1)_m} \frac{x^m}{m!}$$

► So  $A_1 = H \bullet A_2$

$$\begin{aligned} H &= \frac{(\epsilon(2x-1) - x + 1)}{(\epsilon-1)\epsilon(x-1)} \theta + \frac{\epsilon(2x-1) - x + 1}{(\epsilon-1)(x-1)} \\ &= \frac{1}{\epsilon} \theta + \left( 1 - \frac{x}{x-1} \theta \right) + O(\epsilon) \end{aligned}$$



## Step 4 : Differential operator

Example



$$A_1 := \sum_{m=0}^{\infty} \frac{(\epsilon)_m (-\epsilon)_m x^m}{(\epsilon-1)_m m!} \quad , \quad A_2 := \sum_{m=0}^{\infty} \frac{(\epsilon)_m (-\epsilon)_m x^m}{(\epsilon+1)_m m!}$$

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$$A_2 = 1 + \epsilon^2 G(0, 1; x) + \epsilon^3 (-G(0, 0, 1; x) + G(0, 1, 1; x)) + O(\epsilon^4)$$

## Step 4 : Differential operator

### Example



$$A_1 := \sum_{m=0}^{\infty} \frac{(\epsilon)_m (-\epsilon)_m}{(\epsilon-1)_m} \frac{x^m}{m!}, \quad A_2 := \sum_{m=0}^{\infty} \frac{(\epsilon)_m (-\epsilon)_m}{(\epsilon+1)_m} \frac{x^m}{m!}$$

▶ So  $A_1 = H \bullet A_2$

$$\begin{aligned} H &= \frac{(\epsilon(2x-1) - x + 1)}{(\epsilon-1)\epsilon(x-1)} \theta + \frac{\epsilon(2x-1) - x + 1}{(\epsilon-1)(x-1)} \\ &= \frac{1}{\epsilon} \theta + \left(1 - \frac{x}{x-1} \theta\right) + O(\epsilon) \end{aligned}$$



$$A_2 = 1 + \epsilon^2 G(0, 1; x) + \epsilon^3 (-G(0, 0, 1; x) + G(0, 1, 1; x)) + O(\epsilon^4)$$



$$\begin{aligned} A_1 &= 1 + \epsilon \left[ G(1; x) - \frac{x}{x-1} \right] + \epsilon^2 \left[ -\frac{x}{x-1} G(1; x) + G(1, 1; x) - \frac{x}{x-1} \right] \\ &\quad + O(\epsilon^3) \end{aligned}$$

# Mathematica Package

MultiHypExp

MultiHypExp

Available at GitHub [14]

Dependencies :

- ▶ RISC' HolonomicFunctions (Koutschan [15, 16]) : To find the PDE associated with the given MHF and to form the Pfaff system
- ▶ HYPERDIRE (Bytev et.al. [17, 18, 19]) : For step up/down operations
- ▶ CANONICA (Meyer [20]) : To bring the Pfaff system into canonical form
- ▶ PolyLogTools (Duhr et. al. [21]) : To handle MPLs

# Mathematica Package

## MultiHypExp

The package is able to expand the following series

- ▶ **One variable** :  ${}_pF_{p-1}$
- ▶ **Two variables** : Appell  $F_1, F_2, F_3, F_4$ , Horn  $G_1, G_2, G_3, H_1, H_2, H_3, H_4, H_6$  and  $H_7$  and certain KdF functions
- ▶ **Three variables** : Lauricella Saran  $F_A, F_B, F_D, F_K, F_M, F_N$  and  $F_S$
- ▶ Apart from Appell  $F_1, F_2, F_3$  and Horn  $H_2$ , other Appell-Horn series are expanded using their relation to the former functions
- ▶ Series expansion of Appell  $F_4$  and Horn  $H_1$  is possible with certain restriction on the Pochhammer parameters

## MultiHypExp

Commands for one variable

To obtain the series expansion  ${}_2F_1(\epsilon, -\epsilon; \epsilon - 1; x)$

```
In[1]:= SeriesExpand[{{e, -e}, {e - 1}}, {x}, e, 3]
```

```
Out[1]=
```

$$1 + \left(-\frac{x}{-1+x}\right) + G[1, x] e + \left(-\frac{x}{-1+x}\right) - (x G[1, x]) \\ / (-1+x) + G[1, 1, x] e^2 + O[e]^3$$

Alternatively,

```
In[2]:= SeriesExpand[{n}, (Pochhammer[e, n] Pochhammer[-e, n] x^n) \\ / (Pochhammer[e-1, n] n!), {x}, e, 3]
```

yields the same result.

## MultiHypExp

Commands for bi- and tri-variate HF

```
In[3]:= SeriesExpand[F2,{1,1,e,1+e,1-e},{x,y},e,3]
```

```
Out[3]=
```

$$\begin{aligned} & -(1/(-1+x)) + ((-2 G[1,x] + G[1,y] + G[1-y,x]) e) / (-1+x) \\ & + (1/(-1+x)) (2 G[1,x] G[1,y] - 2 G[1,y] G[1-y,x] \\ & + 2 G[0,1,x] + G[0,1,y] - G[0,1-y,x] - 4 G[1,1,x] - 2 G[1,1,y] \\ & + 2 G[1,1-y,x] + 2 G[1-y,1,x] - G[1-y,1-y,x]) e^2 + 0[e]^3 \end{aligned}$$

yields the first three series expansion coefficients of Appell  
 $F_2(1, 1, e; 1 + e, 1 - e; x, y)$  with respect to  $e$  in terms of MPLs.

```
In[4]:= SeriesExpand[{m,n}, exp, {x, y}, e, 3]
```

`exp` must be a series presentation of a MHF with summation indices `m` and `n`.

## MultiHypExp

Commands for obtaining reduction formulae

To find reduction formulae of MHF

```
In[5]:= ReduceFunction[F2, {3, 2, 1, 3, 2}, {x, y}]
```

```
Out[5]=
```

$$1/((-1+x) x (-1+x+y))-G[1,x]/(x^2 y)+G[1-y,x]/(x^2 y)$$

In terms of logarithms

$$F_2(3, 2, 1; 3, 2; x, y) = -\frac{\log(1-x)}{x^2 y} + \frac{\log\left(1 - \frac{x}{1-y}\right)}{x^2 y} + \frac{1}{(x-1)x(x+y-1)}$$

This command can find reduction formulae of Appell  $F_1, F_2, F_3, F_4$  and Lauricella-Saran  $F_D^{(3)}$  and  $F_S^{(3)}$

## MultiHypExp

### Conclusions & Future works

- ▶ Applicable when the parameter  $\epsilon$  appears linearly inside the Pochhammer symbols
- ▶ The package can find the expansion of most of the MHF around integer values of Pochhammer parameters
- ▶ It can find at most first 6 coefficients
- ▶ Not valid at singular locus of the function

**Future works** : Expansion around half-odd integer parameters

- ▶ Some change of variables required but series coefficients can be expressed in terms of MPLs (M.A. Bezuglov et. al [22])
- ▶ May require functions beyond MPLs



Thank You



## References I

- [1] C. Anastasiou, E. W. N. Glover, and C. Oleari. *Nucl. Phys. B* 572 (2000), pp. 307–360. arXiv: hep-ph/9907494.
- [2] F. A. Berends et al. *Z. Phys. C* 63 (1994), pp. 227–234.
- [3] T. Huber and D. Maitre. *Comput. Phys. Commun.* 175 (2006), pp. 122–144. arXiv: hep-ph/0507094.
- [4] T. Huber and D. Maitre. *Comput. Phys. Commun.* 178 (2008), pp. 755–776. arXiv: 0708.2443 [hep-ph].
- [5] S. Moch and P. Uwer. *Comput. Phys. Commun.* 174 (2006), pp. 759–770. arXiv: math-ph/0508008.
- [6] S. Weinzierl. *Comput. Phys. Commun.* 145 (2002), pp. 357–370. arXiv: math-ph/0201011.
- [7] S. Weinzierl. *J. Math. Phys.* 45 (2004), pp. 2656–2673. arXiv: hep-ph/0402131.
- [8] C. Schneider. *Sem.Lothar.Combin.* 56 (2007), pp. 1–36.

## References II

- [9] C. Schneider. “Simplifying Multiple Sums in Difference Fields”. *Computer Algebra in Quantum Field Theory: Integration, Summation and Special Functions*. Ed. by C. Schneider and J. Blumlein. Texts and Monographs in Symbolic Computation. arXiv:1304.4134 [cs.SC]. Springer, 2013, pp. 325–360.
- [10] J. Ablinger. “Computer Algebra Algorithms for Special Functions in Particle Physics”. PhD thesis. RISC, Johannes Kepler University, April 2012.
- [11] Z.-W. Huang and J. Liu. *Comput. Phys. Commun.* 184 (2013), pp. 1973–1980. arXiv: 1209.3971 [physics.comp-ph].
- [12] S. Bera. *Nucl. Phys. B* 989 (2023), p. 116145. arXiv: 2208.01000 [math-ph].
- [13] J. M. Henn. *Phys. Rev. Lett.* 110 (2013), p. 251601. arXiv: 1304.1806 [hep-th].
- [14] S. Bera. *Comput. Phys. Commun.* 297 (2024), p. 109060. arXiv: 2306.11718 [hep-th].
- [15] C. Koutschan. *ACM Commun. Comput. Algebra* 43.3/4 (2010), p. 119.

## References III

- [16] C. Koutschan. *Mathematics in Computer Science* 4.2 (2010), pp. 259–266.
- [17] V. V. Bytev, M. Y. Kalmykov, and B. A. Kniehl. *Comput. Phys. Commun.* 184 (2013), pp. 2332–2342. arXiv: 1105.3565 [math-ph].
- [18] V. V. Bytev, M. Y. Kalmykov, and S.-O. Moch. *Comput. Phys. Commun.* 185 (2014), pp. 3041–3058. arXiv: 1312.5777 [math-ph].
- [19] V. V. Bytev and B. A. Kniehl. *Comput. Phys. Commun.* 189 (2015), pp. 128–154. arXiv: 1309.2806 [math-ph].
- [20] C. Meyer. *Comput. Phys. Commun.* 222 (2018), pp. 295–312. arXiv: 1705.06252 [hep-ph].
- [21] C. Duhr and F. Dulat. *JHEP* 08 (2019), p. 135. arXiv: 1904.07279 [hep-th].
- [22] M. A. Bezuglov and A. I. Onishchenko. (2023). arXiv: 2312.06242 [hep-th].